

The Shape of the Correlation Function

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6 June 2019

- Correlation femtoscopy has become a standard technique for the experimental analysis of heavy-ion collisions
- Two-particle correlation functions are fitted by Gaussian
- However, it seems that the real shape is not Gaussian
- The shape is often better reproduced by Lévy stable distribution
- It has been suggested that non-Gaussian shape may identify the critical point
- We check if the observed shape can be caused by non-critical phenomena

- Correlation function is defined as ratio of two-particle spectrum and one-particle spectra

$$C(p_1, p_2) = \frac{P_2(p_1, p_2)}{P(p_1)P(p_2)} = \frac{E_1 E_2 \frac{d^6 N}{dp_1^3 dp_2^3}}{\left(E_1 \frac{d^3 N}{dp_1^3}\right) \left(E_2 \frac{d^3 N}{dp_2^3}\right)}$$

- We use correlation function in the form

$$C(q, K) - 1 \approx \frac{|\int d^4 x S(x, K) \exp(iqx)|^2}{(\int d^4 x S(x, K))^2}$$

- $K = \frac{1}{2}(p_1 + p_2)$, $q = p_1 - p_2$

Parametrization of correlation function

- Gaussian parametrization

$$C_G(\vec{q}, \vec{K}) = 1 + \lambda \exp \left[- \sum_{i,j=o,s,l} R_{ij}^2 q_i q_j \right]$$

- Lévy parametrization

$$C_L(\vec{q}, \vec{K}) = 1 + \lambda' \exp \left[- \left| \sum_{i,j=o,s,l} R_{ij}'^2 q_i q_j \right|^{\alpha/2} \right]$$

- One-dimensional Lévy parametrisation

$$C_L(Q) = 1 + \lambda' \exp(-|R'Q|^\alpha)$$

- Lévy index characterizes the shape of the correlation function
 - $\alpha = 2 \Rightarrow$ Gauss $\alpha = 1 \Rightarrow$ exponential

Effects leading to non-Gaussianities

- Ensemble averaging
- One-dimensional projection
- \vec{K} averaging
- Resonance decays
- Each event is different
- Averaging over many events may affect the shape

$$C(q, K) \approx 1 + \frac{\langle |\int d^4x S(x, K) \exp(iqx)|^2 \rangle_{ev}}{\langle (\int d^4x S(x, K))^2 \rangle_{ev}}$$

Effects leading to non-Gaussianities

- Ensemble averaging
- Correlation function as a function of a scalar quantity
- One-dimensional projection

- Lorentz-invariant Q

$$Q_{LI}^2 = -q^\mu q_\mu$$

- \vec{K} averaging

- Longitudinally boost-invariant Q

$$Q_{LBI}^2 = \sqrt{(p_{1x} - p_{2x})^2 + (p_{1y} - p_{2y})^2 + q_{long,LCMS}^2}$$

- Resonance decays

Effects leading to non-Gaussianities

- Ensemble averaging
- One-dimensional projection
- \vec{K} averaging
- Resonance decays
- The size of a bin in \vec{K} cannot be taken arbitrarily small
- Correlation function must be averaged over some pair momentum interval

$$C(q, K) \approx 1 + \frac{\int_{bin} d^3 K \left| \int d^4 x S(x, K) \exp(iqx) \right|^2}{\int_{bin} d^3 K \left(\int d^4 x S(x, K) \right)^2}$$

Effects leading to non-Gaussianities

- Ensemble averaging
- One-dimensional projection
- \vec{K} averaging
- Resonance decays
- Different resonances contributes with different lengthscales and timescales
- Correlation function therefore must deviate from a Gaussian form

- This theoretical model is characterized by emission function

$$S(x, p)d^4x = \frac{m_t \cosh(\eta - Y)}{(2\pi)^3} d\eta dx dy \frac{\tau d\tau}{\sqrt{2\pi} \Delta\tau} \exp\left(-\frac{(\tau - \tau_0)^2}{2\Delta\tau^2}\right) \exp\left(-\frac{E^*}{T}\right) \Theta(1 - \bar{r})$$

Cooper-Frye prefactor

- Spatial anisotropy describes shape of the fireball

$$R(\theta) = R_0 \left(1 - \sum_{n=2}^{\infty} a_n \cos(n(\theta - \theta_n)) \right)$$

- Flow anisotropy describes distribution of transverse rapidity

$$\rho(\bar{r}, \theta_b) = \bar{r} \rho_0 \left(1 + \sum_{n=2}^{\infty} 2\rho_n \cos(n(\theta_b - \theta_n)) \right)$$

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Gaussian smearing in proper time

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Boltzmann thermal distribution

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Blast-Wave model

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box profile

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$E^* = p_\mu u^\mu$ - energy in the rest frame

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$\bar{r} = \frac{r}{R(\theta)}$ - scaled radius of fireball in transverse plane

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$$\eta = \frac{1}{2} \ln\left(\frac{t+z}{t-z}\right) \quad \tau = \sqrt{t^2 - z^2}$$

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[Comp. Phys. Comm. 180 (2009) 1642]

- Monte Carlo event generator
- Based on the Blast-Wave model with added resonance decays
- For this study we generated 50,000 events with parameters

$$T = 120 \text{ MeV}$$

$$\tau_{fo} = 10 \text{ fm}/c$$

$$a_2 \in (-0.1; 0.1)$$

$$R = 7 \text{ fm}$$

$$\rho_0 = 0.8$$

$$\rho_2 \in (-0.1; 0.1)$$

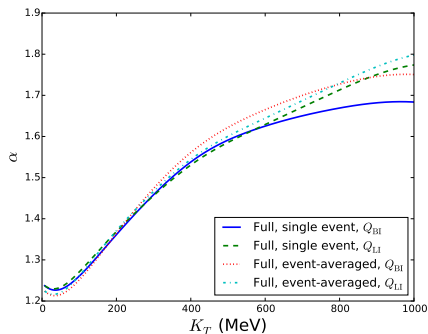
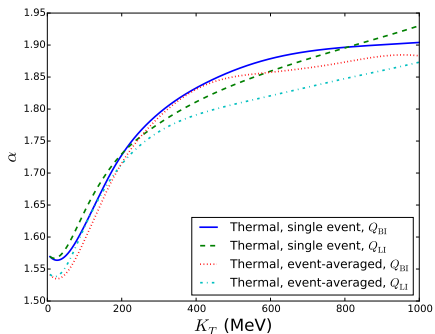
- To generate correlation functions we used **CRAB** [S. Pratt]

- Hydrodynamical model **iEBE-VISHNU** [Comp. Phys. Comm. 199 (2016) 61]
 - 2+1D hydrodynamic simulation
 - Israel-Stewart viscous hydrodynamics
 - Glauber MC initial conditions
- Extension to HBT using **HoTCoffeeh** [Phys. Rev. C 98 (2018) 034910]
- For this study we generated 1,000 events with parameters

$$0 - 10\% \text{ Au+Au collisions at } 200A \text{ GeV}$$
$$T_{fo} = 120 \text{ MeV} \qquad \eta/s = 0.08$$

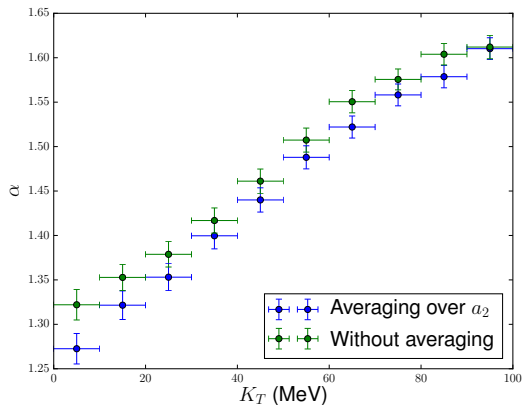
Comparison of different non-Gaussian effects

- Results of hydrodynamical approach
- Comparison of:
thermal vs. full single-event vs. event-averaged Q_{BI} vs. Q_{LI}



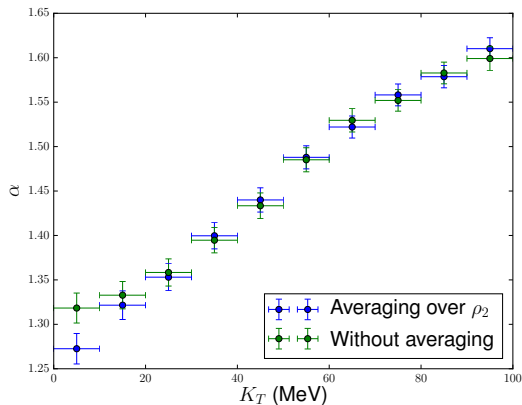
Event-by-event fluctuations

- Study of averaging over events - influence of anisotropies
- Unaveraged - $a_2 = 0.05$ Averaged - $a_2 \in (-0.1; 0.1)$



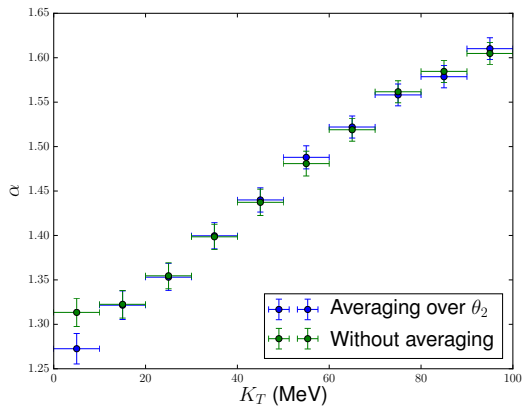
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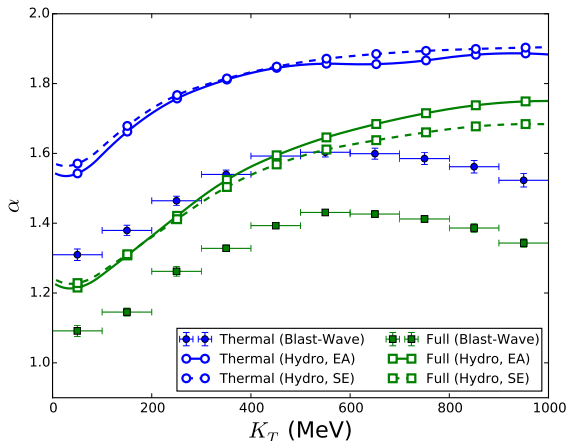


Event-by-event fluctuations

- Study of averaging over events - influence of anisotropies
- Unaveraged - $\theta_2 = 0$ Averaged - $\theta_2 \in (0; \pi)$



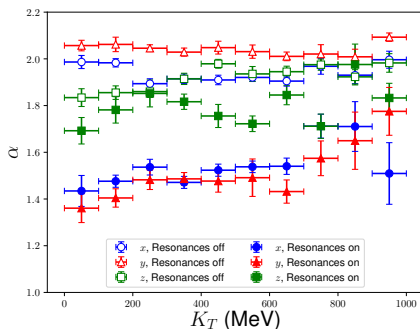
- Resonance decays push down the value of α by 0.2 for both models



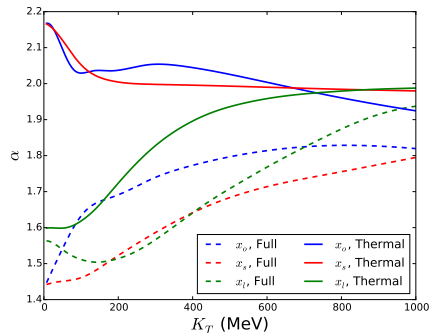
Correlation function in three dimensions

- 3D correlation function fitted via 1D Lévy function in each direction separately

Blast-Wave

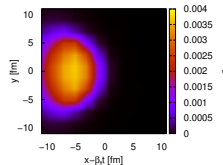
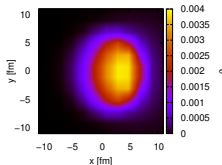
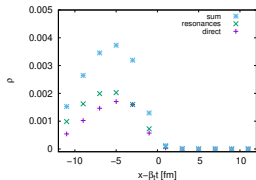
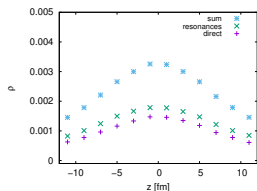
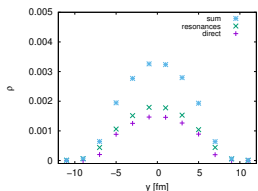
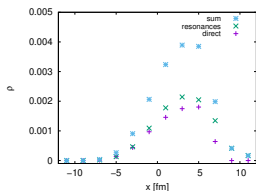


Hydro



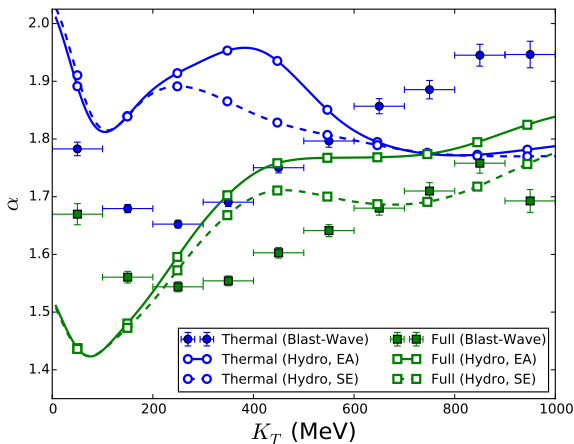
Profiles of emission sources

- To understand the differences in different directions we check the shape of the source which emits pions
- Pions are taken with $K_T \in (300; 400)$ MeV



3D fit to correlation function

- 3D Lévy fit to 3D correlation function



Conclusions

- Lévy index may be influenced by a variety of different mechanisms
- It deviates substantially from the value of 2
- The most significant deviations arise from
 - projection from 3D relative momentum \vec{q} to scalar Q
 - resonance decays
- These conclusions are model-independent